INTERACTION OF A MOIST CAPILLARY-POROUS SOLID WITH A HOT GAS FLOW

G. V. Vasil'eva

UDC 532.63:536.248

An experimental investigation of the temperature distribution in the boundary layer of air above a capillary-porous solid from which a liquid is evaporating is described. An analysis of the experimental results confirms Lykov's hypothesis of the presence of a bulk evaporation source in the boundary layer of air above a capillary-porous solid. This source increases the rate of heat and mass transfer from the solid to the air flow.

Heat and mass transfer between a moist capillary-porous solid and its surroundings is a complicated, interrelated process. Its main feature is that in this case effects taking place within the solid determine the nature of the interaction of the solid and its surroundings. Although such processes are complex, their investigation is of great importance, since similar effects are a feature of such technological processes as drying, evaporative cooling, molecular distillation, etc. In this case evaporation takes place at some depth within the solid and not at its surface. Because of the presence of a dry interlayer of material the heat and mass transfer cannot be calculated from the usual heat transfer formulas. Hence, there is a need to investigate the effect of submersion of the evaporation surface on heat- and mass-transfer coefficients.

This paper gives the results of an experimental investigation of heat and mass transfer from a capillary-porous solid to a gas flow with the evaporation surface inside the solid. The experiments were carried out on the apparatus described in [1]. The test specimen was a model capillary-porous solid made of clear plastic in the form of a $0.2 \times 0.1 \text{ m}^2$ flat plate and filled either with quartz sand or Al₂O₃ powder. The parameters of the external flow in the experiments were $t_{\infty} = 50^{\circ}$ C, $w_{\infty} = 5 \text{ m/sec}$, $\varphi = 7\%$, and the depths of the evaporation front were 0, 2, 4, 6, and 8 mm. We investigated the velocity and temperature distribution in the air boundary layer, the temperature gradient in the dry interlayer of the capillaryporous solid, the temperature gradient on the solid surface in the direction of the flow, and the amount of vapor removed from the solid by evaporation. As was shown in [2], the transverse mass flow has no significant effect on the hydrodynamics of the process. When a liquid evaporates from the pores of a solid a laminar boundary layer is formed at the "solid-air" interface (see Fig. 1). The evaporation of water from the specimen surface was an almost adiabatic process. The process became markedly nonadiabatic when the evaporation front moved into the solid. Significant heat loss from the sides of the experimental solid was prevented by thorough heat insulation. Check measurements showed that heat loss from the sides of the solid and radiative heat transfer between the black walls of the working chamber and experimental solid could be neglected (these two components did not constitute more than 5% of the total heat balance).

Figure 2 shows the temperature distribution in the air boundary layer for quartz sand and alumina powder. The graphs show that the gradient of the curves increases with reduction of the depth of the evaporation surface. At large depths the relationship is almost linear, i.e., it is similar to the case of "pure" heat transfer. The curve has maximum gradient when evaporation takes place at the surface. The temperature of the solid surface increases with increase in the depth of the front and, hence, the temperature difference $t_{\infty} - t_{W}$ is reduced.

The evaporation rate decreases with increase in ξ . A qualitative assessment of the water evaporation rate indicated that evaporation is much more rapid from the pores of Al₂O₃ powder (see Table 1). This fact can be explained in the following way. The quartz sand used had a strictly defined particle size obtained

Institute of Heat and Mass Transfer, Academy of Sciences of the Belorussian SSR, Minsk. Translated from Inzhenerno-Fizicheskii Zhurnal, Vol. 17, No. 3, pp. 463-469, September, 1969. Original article submitted October 24, 1968.

© 1972 Consultants Bureau, a division of Plenum Publishing Corporation, 227 West 17th Street, New York, N. Y. 10011. All rights reserved. This article cannot be reproduced for any purpose whatsoever without permission of the publisher. A copy of this article is available from the publisher for \$15.00.



Fig. 1. Profile of velocity distribution $\omega_X(y)$ in air boundary layer [circles) experiment; dots) calculation].

by passage through chemical sieves. It would presumably have a fairly uniform structure and there would be no great variation in the size of the capillaries. The situation was quite different in the case of Al₂O₃ powder. The mean diameter of the alumina particles were approximately 0.15μ . A powder with such fine pores should have very high porosity (up to 90%). It is known, however, that any disperse material can have a porosity of 46.7%, corresponding to cubic packing, or 26%, corresponding to hexagonal packing. This indicates that in such a capillary-porous substance some Al₂O₃ particles form solid inclusions. This produces a highly inhomogeneous system with capillaries (r < 10^{-5} cm) and large cavities (this can be seen by using a microscope with a binocular attachment). In a capillary-porous solid of this structure the moisture distribution is altered. The capillaries extract water from the large pores and the pressure drop in the dry interlayer will depend considerably on its capillary-porous structure. This leads to an increase in the flow rate of the evaporated liquid.

We will give an account of the method of calculating the heat fluxes and heat-transfer coefficient. The amount of heat passing from the hot air flow into the capillary-porous solid was calculated

in two ways. The first method involved the temperature profile in the air boundary layer

$$q = -\lambda_a \left(\frac{\partial t}{\partial y}\right)_{y \to 0};\tag{1}$$

while the second method used the amount of heat passing through the dry interlayer of the capillary-porous solid due to its heat conduction

$$q = -\lambda_i \frac{t_w - t_{\xi}}{\xi} \,. \tag{2}$$

The good agreement of the results obtained by the two methods indicates that they are equally applicable to the described process (see Table 2).

It is of interest to note that the method of heat flux calculation used in the case of evaporation of liquid from a solid surface

$$q = jr + q_{\rm h} \tag{3}$$

(in adiabatic conditions the second term is omitted from the equation) is quite unsuitable in the case of evaporation of a liquid from within a capillary-porous solid. The heat which can pass through the dry interlayer of the solid is not sufficient for evaporation of the amount of water which is converted to vapor. A saturated vapor-air mixture moves through the pores in the solid. Since the matrix of a capillary-porous



Fig. 2. Temperature distribution over quartz sand (a) and over Al_2O_3 powder (b) (y, mm; t, °K): 1,2,3,4,5) thickness of dry interlayer 8,6,4,2, and 0 mm, respectively.

	٤. mm							
Material	0	2	4	6	8			
Quartz sand Al ₂ O3 powder	2,1	1,1 2,2	0,85 2,1	0,70 1,8	0,52 1,2			

TABLE 1. Flow Rate of Evaporated Water j (kg/m² \cdot h)



Fig. 3. Plate surface temperature distribution along x-axis (t, $^{\circ}$ K; *l*, m): 1, 2, 3, 4, 5) thickness of dry interlayer 0, 2, 4, 6, and 8 mm, respectively.

solid offers a fairly high hydraulic resistance to mass transfer, there will be no flow of moisture unless there is a pressure difference in the pores. The greater the thickness of the dry interlayer, the greater the pressure gradient. The moist vapor air mixture will be expelled and will subsequently evaporate in the air flow over the plate surface. Hence, we can postulate the presence of an additional heat sink ("bulk" evaporation) over a capillary-porous solid.

Table 2 gives the values of the heat-transfer coefficient between the solid surface and the gas flow.

The tabulated data indicate that the heat transfer coefficient increases with increase in the depth of the evaporation front. This is due to the fact that an increase in ξ causes an increase in the temperature of the solid surface and a reduction of the heat flux. The changes

in q and t_w are not directly related. Both the heat flux and the solid surface temperature depend on the thermal resistance of the dry interlayer and the thermal resistance of the air boundary layer. Because of the complex structure of the capillary-porous solid in which evaporation occurs the heat flux decreases more slowly than the surface temperature increases, which leads to an increase in the heat-transfer coefficient. This again confirms the fact that a rigorous formulation of the problem of heat transfer between a damp capillary-porous solid and its surroundings will necessitate the joint consideration of the heat transfer transfer inside the solid and the interrelated heat and mass transfer in the boundary layer close to its surface. Neglect of particular factors (flow temperature, velocity, humidity, porousity of solid, thermal conductivity of solid, thickness of dry interlayer) can lead to quantitatively and qualitatively contradictory results.

To find out if the heat-transfer coefficient is increased or reduced when liquid evaporates from a solid in comparison with "pure" heat transfer, we experimentally measured the temperature profile in the air boundary layer and determined the heat-transfer coefficient, uncomplicated by mass transfer, by the method described above. The ratio α_{ξ}/α_0 or Nu_{ξ}/Nu_0 , denoted by the dimensionless complex N, was greater than unity. This indicates that in the case of low-rate processes the transverse mass flow leads to an increase in the heat-transfer coefficient. This was confirmed by the temperature distribution along the solid surface. In the case of adiabatic evaporation t_W will be constant and equal to the temperature of a wet-bulb thermometer in the flow, and the temperature head Δt will be $t_{\infty} - t_W$. If the process is non-adiabatic the solid surface temperature will differ from the reading of a wet-bulb thermometer. Figure 3 shows the solid surface temperature profile along the flow for quartz sand and an external flow with $t_{\infty} = 50^{\circ}C$ and $w_{\infty} = 5 \text{ m/sec}$. Figure 3 shows that the temperature head increases in the direction of the gas flow. Since the flow velocity and temperature differences in the experiments were small the temperature head along the plate was 1-3^{\circ}C. A higher temperature head can be expected at higher temperatures and velocities.

It was also noted that with increase in the interlayer thickness the temperature difference between the beginning and end of the plate increased. It is known from heat-transfer theory that if the temperature head in the direction of the heat flow on the surface of a solid is increased, the heat-transfer coefficient is greater than at constant surface temperature. Hence, with increase in the thickness of the dry interlayer the heat-transfer coefficient between the flow and the solid surface will increase.

-											
٤	<i>q</i> ₁	q 2	<i>q</i> 3	α	Nu	$N = \frac{Nug}{Nu_0}$	Н	Ķ	N*		
Quartz sand											
$ \begin{array}{c} 0 \\ 2 \\ 4 \\ 6 \\ 8 \end{array} $	212 156 122,9 78,0 62,5	202 148 117,5 70,8 57,5	1180 985 890 602 553	10,1 10,34 11,0 11,28 11,5	72 74 79,2 80,5 82,1	1,23 1,27 1,35 1,37 1,4	524 262 175 131	0,427 0,2135 0,1428 0,107	1,26 1,32 1,40 1,46		
Alumina powder											
	162 150 106 70,8	156 135 119 102	1627 1590 1440 1057	9,32 10,2 10,5 10,9	66,7 73,2 76,3 78,5	1,14 1,25 1,30 1,33	1430 715 477 358	1,165 0,582 0,389 0,292	1,15 1,20 1,25 1,30		

TABLE 2. Heat Characteristics in Relation to Thickness of Dry Interlayer of Material

Note: N* is calculated from Eq. (3).

The obtained experimental data were also correlated by general relationships obtained from a solution of the differential transport equation in a boundary layer with allowance for the effect of submersion of the evaporation surface [2, 3].

The dimensionless variable $K = Hx/\sqrt{Pe}$ (where $H = \lambda_a/\lambda_S\xi$) characterizes the effect of submersion of the evaporation surface on the heat-transfer coefficient; $B = (w_y/w_x)Re_x^{0.5}Pr^{0.5}$ is a dimensionless parameter characterizing mass transfer in the transverse direction. If we assume a parabolic velocity distribution $w_x(y)$ along the plate on the y axis, then B = 0.21 (see [3]). The quantity N characterizes the relative increase in the local Nusselt Number due to evaporation of water from a capillary-porous solid

$$N(K, B) = \frac{\sqrt{\pi N u_x}}{1 \ \overline{Pe_x}} = \frac{\varphi(K, B) - \frac{1}{2}}{\left(1 - \frac{B}{K}\right) - \frac{1}{K\sqrt{\pi}}\varphi(K, B) + \frac{1}{2}} \frac{B}{K} \operatorname{erfc} \frac{1}{2}B}, \qquad (4)$$

where

$$\varphi(K, B) = \left(1 - \frac{1}{2} \quad \frac{B}{K}\right) \cdot \overline{\pi} K \exp(K^2 - BK) \operatorname{erfc}\left(K - \frac{1}{2} B\right).$$
(5)

The values of H and K were determined for quartz sand and Al_2O_3 powder and the values of N were calculated from formulas (4) and (5). Table 2 gives the results of the calculations, which also confirm the conclusion that the heat-transfer coefficient is higher in the case of evaporation from capillary-porous solids than in the case of a dry solid.

The experimental measurements of the thermal boundary layer in the case of evaporation of water from Al_2O_3 powder and quartz sand into a flow of hot air lead to the following conclusions:

- 1. The thermal boundary layer associated with evaporation of a liquid depends on the depth of the evaporation surface.
- 2. The heat flux on the wall $q = -\lambda_a (\partial t/\partial y)_{y \to 0}$ is many times greater than the heat flux given by Eq. (3), which indicates that vapor-gas mixture is expelled from the pores and subsequently evaporates in the air flow and its boundary layer.
- 3. Evaporation of water from the pores of alumina powder is more rapid than from quartz sand. This can be attributed to the complex structure of Al_2O_3 powder. Its high porosity (about 90%) is due to the fact that there are large cavities between particle conglomerates, in addition to micropores ($r < 10^{-5}$ cm).
- 4. When a dry interlayer is present the heat flux depends on the thermal resistance of the boundary layer and the dry interlayer. The surface temperature increases and the heat flux decreases with increase in ξ . There is no direct relation between the change in these quantities. In view of the complex structure of the capillary-porous solid in which evaporation occurs the heat flux decreases

more slowly than the surface temperature increases, and, hence, since $q = \alpha (t_{\infty} - t_W)$, there is an increase in the rate of emission of vapor-water mixture from the pores with increase in ξ and, hence, an increase in the heat-transfer coefficient.

NOTATION

- t is the temperature;
- w is the velocity;
- φ is the moisture content;
- ξ is the thickness of dry interlayer of solid;
- r is the pore radius;
- q is the specific heat flux;
- x, y are the coordinates;
- λ is the thermal conductivity;
- j is the mass flow density;
- α is the heat-transfer coefficient;
- *a* is the thermal diffusivity;
- Nu is the Nusselt number;
- Pr is the Prandtl number;
- Re_x is the local Reynolds number;
- Pe_X is the local Peclet number referred to mean integral velocity in boundary layer $(Pe_X = w_X x/a)$.

Subscripts

- ∞ denotes the flow;
- w denotes the wall;
- 0 denotes the dry solid.

LITERATURE CITED

- 1. G. V. Vasil'eva, Inzh.-Fiz. Zh., 9, No. 3 (1965).
- 2. A. V. Lykov, Drying Theory [in Russian], Énergiya (1968).
- 3. A. V. Lykov and G. V. Vasil'eva, Inzh.- Fiz. Zh., 14, No. 3 (1968).